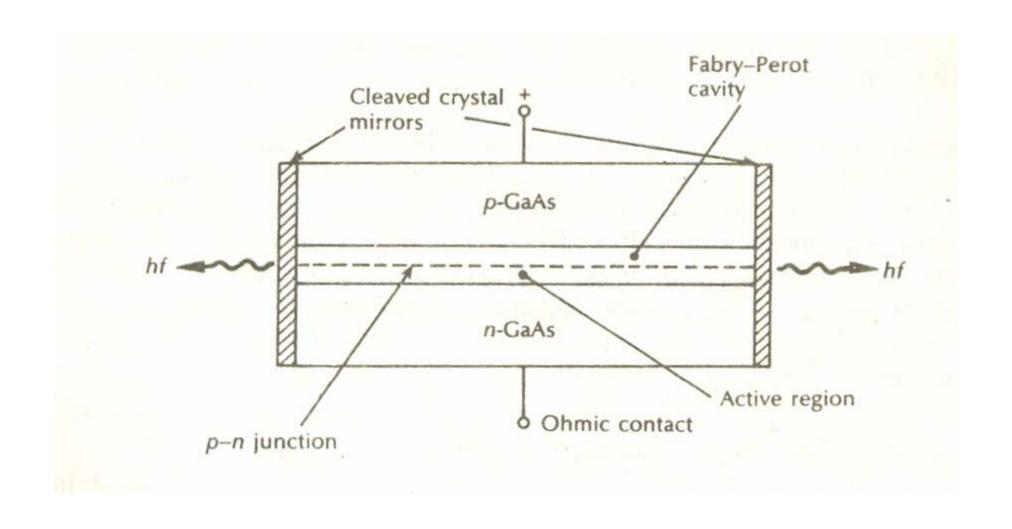
# LASERS

#### **SEMI CONDUCTOR INJECTION LASER DIODE**

<u>Principle</u>: Stimulated emission by recombination of injected carriers by provision of an optical cavity to provide the feed back of photons

## Advantages of ILD:

- 1. High radiance ( due to amplifying effect of stimulated emission).
- 2. Narrow line width (1 nm) & so less dispersion.
- 3. Better modulation capabilities.
- 4. More coherent & hence better focusing capability.
- Better coupling efficiency of output power into optical fiber



GaAs homojunction injection laser with a Fabry – Perot cavity

### EINSTEIN RELATIONS-LASER

 Atomic system is in thermal equilibrium, i.e. the rate of upward transitions must be equal to rate of downward transitions

 $N_1$ ,  $N_2$  = density of atoms in energy levels  $E_1$  &  $E_2$   $g_{1,}$   $g_2$  = corresponding degeneracy's (no. of sub levels within energy levels  $E_1$ ,  $E_2$ )

$$N_1 / N_2 = g_1 e^{-E_1/KT}/g_2 e^{-E_2/KT}$$
  
=  $g_1 / g_2 e^{(E_2 - E_1)/KT} = g_1/g_2 e^{hf/KT}$  (1)

Where,  $E_2 - E_1 = hf = energy of single photon K- Boltzmann's constant, T= abs. temp.$ 

 $\rho f$  – spectral density,  $B_{12}$  – Einst. coefft. of absorption

$$R_{12}$$
 = rate of upward transition =  $N_1 \rho f B_{12}$   
 $(R_{12} \alpha N_1 \rho f)$ 

R<sub>21</sub> = downward transition rate = sum of spontaneous and stimulated contributions

$$R_{21} = N_2 A_{21} + N_2 \rho f B_{21}$$

where  $A_{21}$  = Einst. Cofft.of spontaneous emission =  $1/t_{21}$ 

 $t_{21}$  = spontaneous life time

For thermal equilibrium,  $R_{12} = R_{21}$ 

$$N_1 \rho f B_{12} = N_2 A_{21} + N_2 \rho f B_{21}$$

$$\rho f (N_1 B_{12} - N_2 B_{21}) = N_2 A_{21}$$

$$\rho f = N_2 A_{21} / N_1 B_{12} - N_2 B_{21} = (A_{21} / B_{21}) / (N_1 B_{12} / N_2 B_{21}) - 1$$

Putting 
$$N_1/N_2 = g_1/g_2 \, e^{\, hf/Kt} \, from = n(1)$$
  $\rho f = (A_{21}/\,B_{21})/(g_1\,B_{12}/g_2\,B_{21}) \, e^{\, hf/Kt} \, -1$   $\rho f = (A_{21}/\,B_{21})/[\,\,(g_1B_{12}/g_2\,B_{21})e^{\, hf/Kt}\,\,] \, -1 = radiation$  density from a black body (2) Radiation density for a black body (Planck's Result )  $\rho f = 8\pi \, hf^{\,3}/c^3(1/\,e^{\, hf/\,KT} - 1)$  (3) Comparing the above two results of eq. 2 & 3, we have  $-(A_{21}/B_{21})/(g_1B_{12}/g_2\,B_{21})e^{\, hf/KT} - 1 = 8\pi \, hf^{\,3}/c^3\,(1/\,e^{\, hf/\,KT} - 1)$   $A_{21}/B_{21} = 8\,\pi \, hf^{\,3}/\,c^3\,\,(B)$   $(g_1B_{12}/g_2\,B_{21})e^{\, hf/KT} - 1 = e^{\, hf/Kt} - 1$   $(g_1B_{12}/g_2\,B_{21}) = 1 \, \text{or} \, B_{12} = (g_2/g_1)\, B_{21}$ 

If 
$$g_1 = g_2$$
,  $B_{12} = B_{21}$ 

Einstein Coefft. of absorption = Einstein coefft of stimulated emission

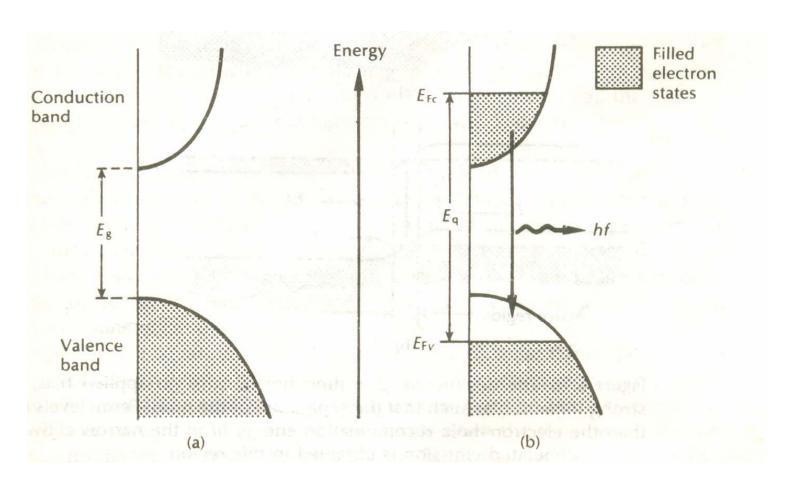
From = n A and = n B

$$\rho f = A_{21}/B_{21} (1/e^{hf/KT} - 1)$$

Or 
$$B_{21} \rho f / A_{21} = (1/e^{hf/KT} - 1)$$

= stimulated emission rate / spontaneous emission rate

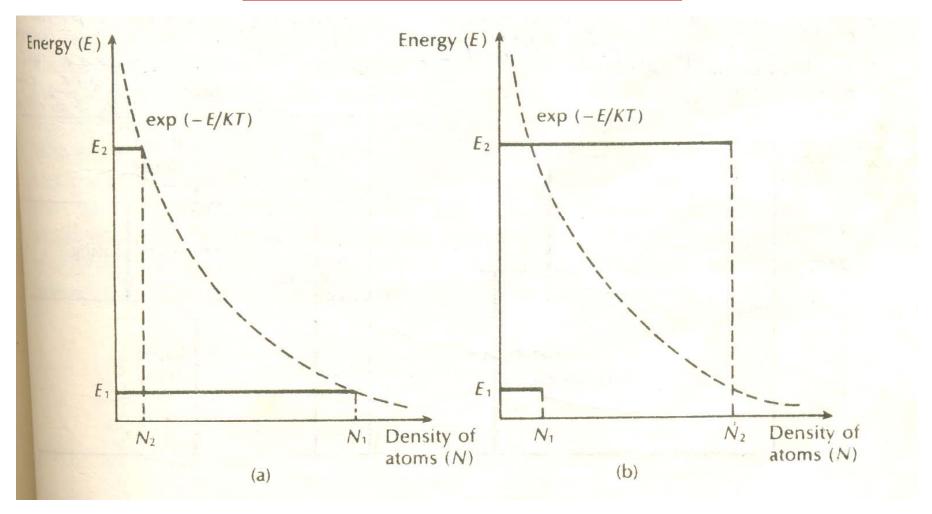
## **STIMULATED EMISSION & LASING:-**



The filled electron state for an intrinsic direct band gap semiconductor at absolute zero (a) in equilibrium (b) with higher carrier injection

- At Abs. zero temperature, conduction band contains no electrons.
- Incident photons with energy Eg < ( $E_{fc} E_{fv}$ ) can't be absorbed
- These photons induce downward transition of an electron and thus stimulated emission of another photon.
- Basic condition for stimulation :  $E_{fc} E_{fv} > hf > Eg$
- Heavy doping satisfies the above condition in a p-n junction thereby providing stimulated emission whereas in a normal p-n junction only spontaneous emission takes place (LED).
- Another condition (for pn diode) to establish lasing is the <u>provision of optical feedback to give laser</u> <u>oscillation (Fabry – Perot Cavity)</u>

## **POPULATION INVERSION**

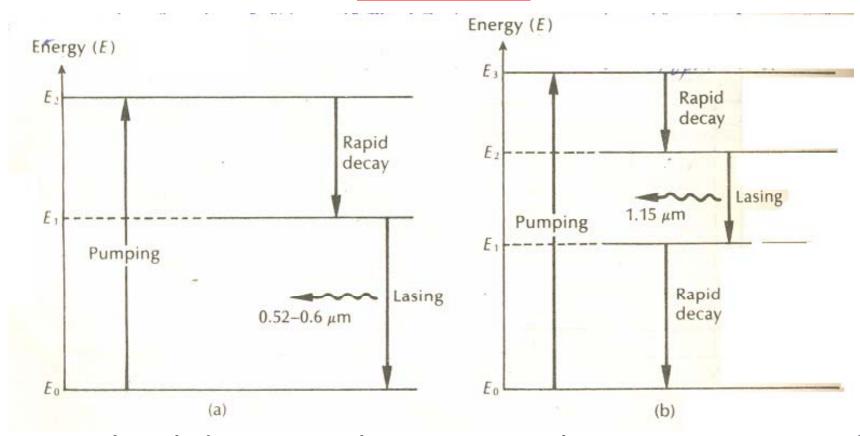


(a)Boltzmann distribution for a system in thermal equilibrium; (b) A non-equilibrium distribution showing population inversion

# **FOR OPTICAL AMPLICATION**: N<sub>2</sub> > N<sub>1</sub> (to increase the rate of stimulated emission)

- <u>PUMPING</u>: An external energy source is used to excite atoms into the upper energy state, thereby obtaining condition of population inversion.
- Intense radiation from a HF radio field or an optical flash tube is applied.
- The two level system does not lend itself to suitable population inversion.
- From = n  $B_{12} = (g_2/g_1) B_{21}$ At best  $B_{12} = B_{21}$  (when,  $g_1 = g_2$ ) which means prob. of absorption and stimulated emission are equal  $(N_1=N_2)$
- Hence, the necessity to use three or four energy level to obtain population inversion.

## THREE AND FOUR LEVEL SYSTEMS - POPULATION INVERSION



Energy level diagrams showing population inversion and lasing for two non-semiconductor lasers: (a) three level system — ruby (crystal) laser; (b) four level system — He—Ne (gas) laser.

- Both systems display a central metastable state in which atoms spend an usually long time.
- Stimulated emission (lasing) takes place from this metastable state.
- With pumping, electrons move from  $E_0$  to  $E_2$  state.
- Electrons rapidly decay by non- radiative process to either  $E_1$  or  $E_0$ , providing empty states in  $E_2$ .
- Thus density of atoms in metastable state  $N_1$  increases above the ground state  $N_0$ .
- $N_1 > N_0$  (population inversion) LASING

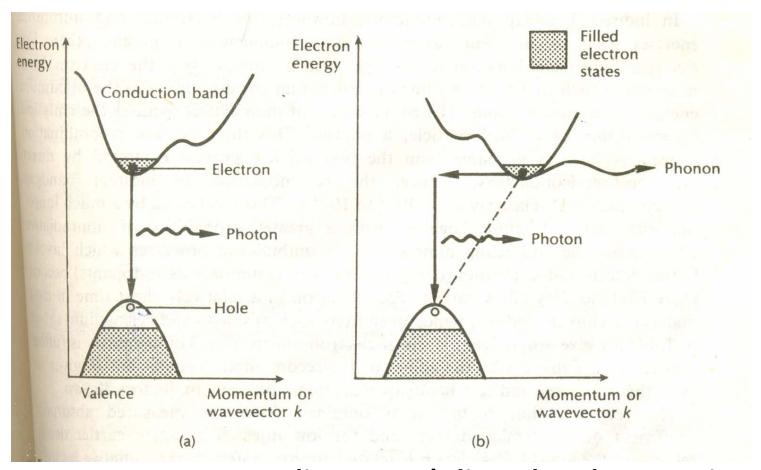
## Three level system requires high pump powers

## Four level system

Pumping excites atoms from E<sub>0</sub> to E<sub>3</sub>, and they decay rapidly to meta stable level E<sub>2</sub>

Since population of  $E_3$  and  $E_1$  remain essentially unchanged a small increase in no of atoms in  $E_2$  creates population inversion hence lasing takes place between  $E_2$  &  $E_1$ 

 FOUR LEVEL SYSTEM HAS MUCH LOWER PUMPING REQUIREMENTS



Energy — momentum diagrams a) direct bandgap semiconductor; (b) indirect bandgap semiconductor.

Recombination is slow in indirect BG semiconductor

Thus competing non-radiative recombination process becomes more likely (due to impurities/ defects)

DIRECT & INDIRECT BANDGAP SEMICONDUCTORS

## Therefore indirect band gap materials give insignificant levels of electroluminescence

 $p = 2\pi hk$ 

Avg minority carrier lifetime= $10^{-8}$  to  $10^{-10}$  sec(DBGSC) =  $10^{-2}$  to  $10^{-4}$  sec(IDBGSC)

Material Energy Band gap (eV) Recom.coefft.

GaAs Direct: 1.43 7.21×10<sup>-10</sup>

GaSb Direct: 0.73  $2.39 \times 10^{-10}$ 

InAs Direct: 0.35 8.5×10<sup>-11</sup>

InSb Direct: 0.18  $4.58 \times 10^{-11}$ 

Si Indirect: 1.12  $1.79 \times 10^{-15}$ 

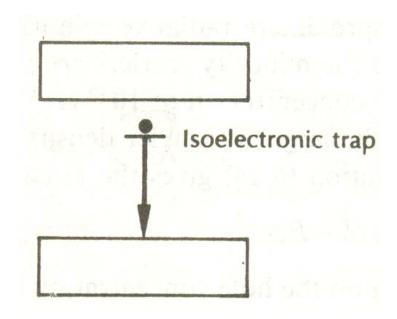
Ge Indirect: 0.67  $5.25 \times 10^{-14}$ 

GaP indirect: 2.26  $5.37 \times 10^{-14}$ 

## Addition of impurity center (N<sub>2</sub>) to GaP, indirect BG material converts it into a direct BG material

#### **Nitrogen:-ISOELECTRONICIMPURITY**

(As it has the same no of valance electrons as Phosphorus)



Nitrogen impurity center captures an electron and acts as an ISOELECTRONIC trap which has a large spread of momentum.

## **Nitrogen:-ISOELECTRONICIMPURITY**

This trap attracts the oppositely charged carrier (hole) and a direct transition takes place between impurity centre and valence band. GaP + N<sub>2</sub> becomes an efficient light emitter. The carrier lifetime is effectively reduced.

## **SINGLE MODE OPERATION:**

The spectral width of the emission from the single mode device is far smaller than the broadened transition line width, which occurs due to

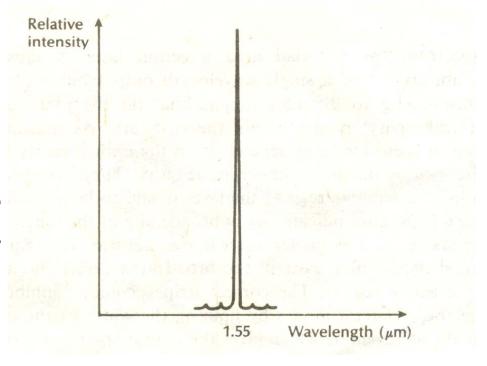
 a) emission over a small frequency band instead on a single frequency.

b) frequency variations due to thermal motion of atoms -

Doppler broadening) )

c) due to atomic collisions

Typical single longitudinal mode output spectrum from a single-mode injection laser.



## **OBTAINING A SINGLE MODE OPERATION.**

This can be done by reducing the length of the cavity, L until frequency of separation  $\delta f = c/2nL$  is larger than the laser transition line width or gain curve. Then only the single mode which falls within the transition line width can oscillate within the laser cavity.

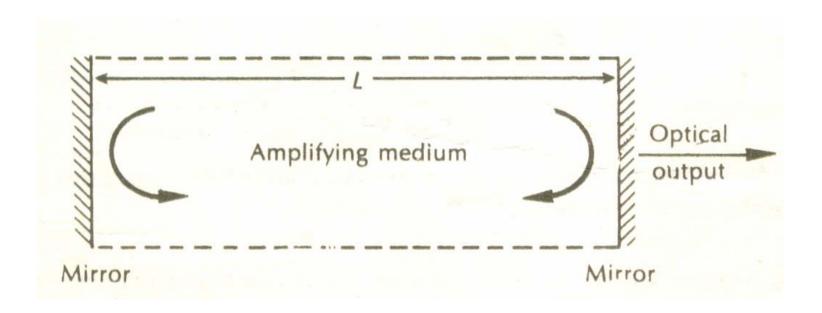
Note: Injection lasers with short cavity lengths ( =50  $\mu$ m) are difficult to handle and have not been very successful.

## **OPTICAL FEEDBACK AND LASER OSCILLATION**

Light amplification in laser occurs due to a photon colliding with an atom (in the excited energy state) to cause stimulated emission of a second photon. Then both these photons release two more.

CONTINUATION OF THIS PROCESS CREATES AVALANCHE MULTIPLICATION.

When em waves associated with these photons are in phase, amplified coherent emission is obtained.



The basic laser structure incorporating plane mirrors

The optical cavity formed provides positive feedback of photons by reflection at the mirrors. The structure acts as a FABRY – PEROT resonator.

One mirror is made partially transmitting for useful radiation to escape from cavity.

## THRESHOLD CONDITION FOR LASER OSCILLATION

## CONDITIONS FOR LASER OSCILLATION

- Population inversion
- Threshold gain within the amplifying medium ( to initiate and sustain the laser oscillation).

Let  $\alpha$  = loss coefficient per unit length.

 $r_1$ ,  $r_2$  = reflectivity of two mirrors.

Fractional loss incurred on each round trip of the beam =  $r_1 r_2 e^{-2 \alpha L}$ 

# Fractional gain = $e^{2gL}$ (due to stimulated emission) where g = gain coefficient\_per unit length Hence $e^{2gL} \times r_1 r_2 e^{-2\alpha L} = 1$

$$\begin{split} r_1 r_2 &\, e^{\, 2L(g-\, \alpha)} = 1 \\ \text{Or } e^{\, 2L(g-\, \alpha)} &= \, 1/r_1 r_2 \quad \text{or } \, 2L\, (g-\alpha) = \log\, 1/r_1 r_2 \\ (g-\alpha) &= \, 1/2L[\log\, 1/r_1 r_2] \\ g &= \, l/2L\, [\log\, 1/r_1 r_2] + \alpha \\ (\, trans.\, loss\, thr'\, mirrors\, ) \end{split}$$

## Assignment

- What is semiconductor laser diode? Give its advantages and disadvantages.
- Derive EINSTEIN relation for lasers.
- What is Threshold condition for laser oscillation?
- Explain three and four level system population inversion.